

Anisotropy-driven enhancement of spin correlations

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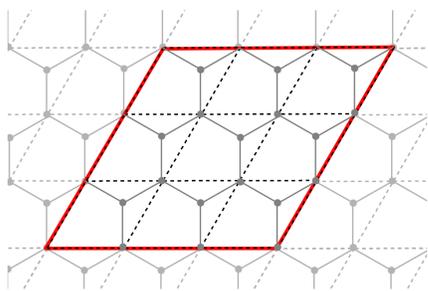
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We study the Hubbard model on different lattices - coupled 1D chains, coupled 2D layers made of square lattice, layered honeycomb lattice - and investigate the thermodynamic properties by the dynamical cluster approximation. We find that the short-range spin correlations are significantly enhanced for the anisotropic models in the direction with stronger tunneling amplitudes when compared to the isotropic 3D cubic system. Our results provide a thermometer for the quantum simulation experiment of ultracold fermions in an optical lattice and allow an quantitative estimate of the excess entropy during the lattice loading. We furthermore investigate the dependence of the critical temperature (entropy) at the Néel transition on anisotropy and lattice geometry.

DCA METHOD

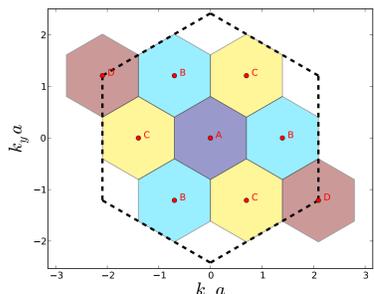
The method of choice is the dynamical cluster approximation (DCA) [2], a cluster extension of the dynamical mean-field theory (DMFT) [1]. The lattice problem is self-consistently mapped onto an impurity (=cluster) with periodic boundary conditions and additional non-interacting bath.



Cluster with 9 unit cells of the hexagonal lattice in the real space.

The DMFT provides exact solution of the Hubbard model in an infinite dimensional system - in that case the self-energy is purely local (\rightarrow i.e. \mathbf{k} -independent). That motivates the (only) approximation in any DMFT method: equality of the impurity and lattice self-energy - in particular for DCA we require

$$\Sigma^{\text{imp}}(\mathbf{K}) \approx \Sigma^{\text{lat}}(\mathbf{k}) \quad (\text{constant over patch}).$$



9-unit-cells cluster on the hexagonal lattice in reciprocal space, with hexagonal patches.

The mapping is determined by the self-consistency condition

$$G^{\text{imp}}(\mathbf{K}) = \frac{1}{\Omega} \int_{\mathbf{K}\text{-patch}} d\mathbf{k} G^{\text{lat}}(\mathbf{k}).$$

The impurity solver for the DCA calculations used for this work was the continuous time auxiliary field quantum Monte Carlo impurity solver (CT-AUX) [3] [4].

$$\begin{aligned} & \text{Monte Carlo solver: } G_{\text{imp}}^0(\mathbf{K}) \\ G_{\text{imp}}^0(\mathbf{K}) &= [\bar{G}(\mathbf{K})^{-1} + \Sigma(\mathbf{K})]^{-1} & \Sigma(\mathbf{K}) &= [G_{\text{imp}}^0(\mathbf{K})^{-1} - G_{\text{imp}}(\mathbf{K})^{-1}] \\ \bar{G}(\mathbf{K}) &= \int_{\mathbf{K}\text{-patch}} [G^0(\mathbf{k})^{-1} - \Sigma(\mathbf{K})]^{-1} \end{aligned}$$

Iterative solution scheme

The DCA method is a controllable approximation and it becomes exact in the limit of infinite cluster size. The mean-field character remains in the vicinity of the phase transitions of 2nd kind, where the correlation length is larger than the cluster size and thus it cannot be captured by the impurity. Away from the phase transitions the observables may be extrapolated to the thermodynamic limit using the known asymptotical behavior.

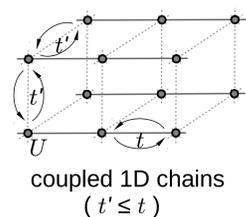
A way to measure the susceptibility and to find the phase transition is usage of the pairing matrix formalism with an additional approximation on the vertex function,

$$\Gamma^{\text{imp}}(\mathbf{Q}m, \mathbf{K}n, \mathbf{K}'n') \approx \Gamma^{\text{lat}}(\mathbf{Q}m, \mathbf{k}n, \mathbf{k}'n').$$

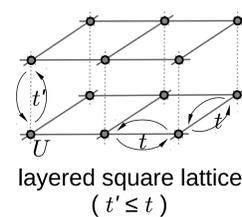
MODEL

We study the anisotropic Hubbard model on 3-dimensional lattices described by the Hamiltonian

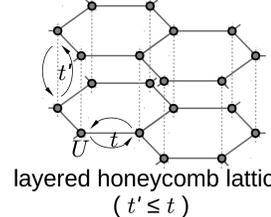
$$\hat{H} = - \sum_{\langle i,j \rangle, \sigma} t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow} - \mu \sum_{i,\sigma} \hat{n}_{i\sigma}.$$



coupled 1D chains ($t' \leq t$)

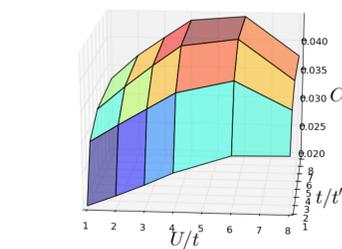


layered square lattice ($t' \leq t$)



layered honeycomb lattice ($t' \leq t$)

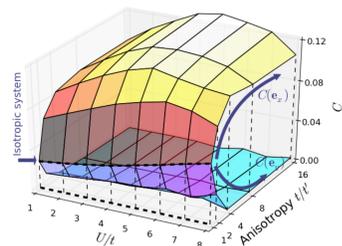
SHORT-RANGE SPIN-SPIN CORRELATIONS $C = -2 \langle S_{\mathbf{r}}^z S_{\mathbf{r}+\Delta}^z \rangle$



Layered lattices at $S/N = 0.8$ at half filling

\leftarrow layered square

layered honeycomb \rightarrow



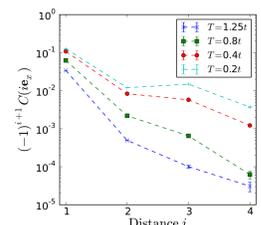
Coupled 1D chains at half filling [6]

\leftarrow at temperature $T = t/2$

Upper sheet: in-chain direction

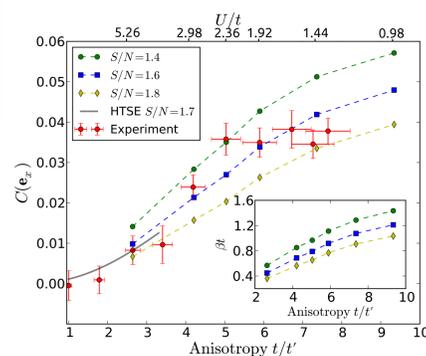
Lower sheet: \perp direction

$$U = 1.44t, t/t' = 7.36 \rightarrow$$



Enhancement of nearest-neighbor spin-spin correlation roughly proportional to the inverse coordination number ratio, counting hoppings t only.

COMPARISON WITH EXPERIMENT [5] UTILIZING LDA



Coupled 1D chains: comparison with experiment [6]

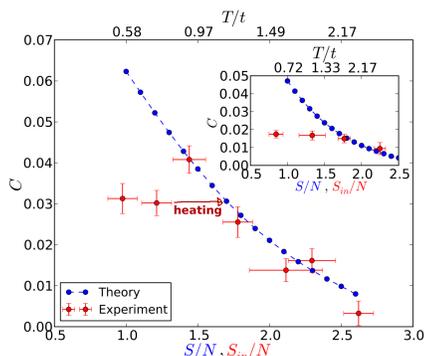
\leftarrow Thermometer: $C \rightarrow T/t$

Estimate of heating \rightarrow

during the lattice loading

Main: $U = 1.44t, t/t' = 7.36$

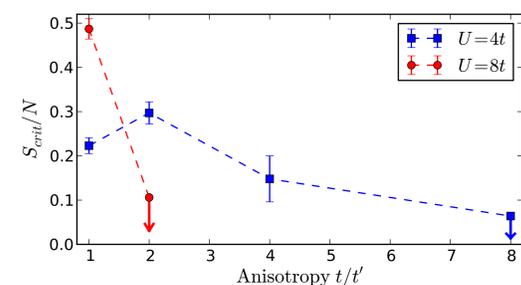
Inset: $U = 2.98t, t/t' = 4.21$



LONG-RANGE ANTIFERROMAGNETIC ORDER AT HALF FILLING

lattice	t/t'	U/t	S_{crit}/N
layered square	2.0	6.0	0.45(3)
layered square	2.0	8.0	0.37(12)
layered honeycomb	1.0	8.0	< 0.49

coupled 1D chains [6] \rightarrow



Cubic lattice is in terms of highest S_{crit}/N optimal.

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